Dissipation enhancement for stochastic heat equation with transport noise

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Based on arXiv:2104.01740

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16th Workshop on Markov Processes and Related Topics

BNU & CSU, 2021.7.12

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Stabilization by noise

• For a $d \times d$ matrix A_0 , the trivial solution $X_t \equiv 0$ of ODE

$$\dot{X}_t = A_0 X_t, \quad X_0 \in \mathbb{R}^d$$

is unstable if A_0 has positive eigenvalues.

• When can the above system be stabilized by noise?

$$dX_t = A_0 X_t dt + \sum_{i=1}^m A_i X_t \circ dW_t^i, \quad X_0 \in \mathbb{R}^d.$$

Here A_1, \ldots, A_m are some matrices.

• Define the Lyapunov exponent

$$\lambda(X_0, \omega) = \limsup_{t \to \infty} \frac{1}{t} \log |X_t(X_0, \omega)|.$$

Oseledec's multiplicative ergodic theorem: with prob. 1, $\exists d$ random numbers $\lambda_1 \leq \cdots \leq \lambda_d$ such that $\lambda(X_0, \omega)$ takes one of them.

Stabilization by noise

• L. Arnold-Crauel-Wihstutz (SIAM J Contr. Optim., 1983): for some $\nu > 0$, consider

$$dX_t = A_0 X_t dt + \nu \sum_{i=1}^m A_i X_t \circ dW_t^i, \quad X_0 \in \mathbb{R}^d.$$

 ν is the intensity of the noise.

Theorem

There are skew-symmetric matrices A_1, \ldots, A_m such that the top Lyapunov exponent λ_{ν} satisfies

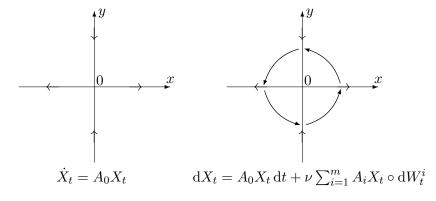
$$\lim_{\nu \to \infty} \lambda_{\nu} = \frac{\operatorname{Tr}(A_0)}{d}.$$

In particular, $\dot{X}_t = A_0 X_t$ can be stabilized by noise iff $\text{Tr}(A_0) < 0$.

$$d|X_t|^2 = 2\langle X_t, A_0 X_t \rangle dt + 2\nu \sum_{i=1}^m \underbrace{\langle X_t, A_i X_t \rangle}_{=0} \circ dW_t^i = 2\langle X_t, A_0 X_t \rangle dt.$$

A trivial example

Consider
$$A_0 = \begin{pmatrix} 1 & 0 \\ 0 & -2 \end{pmatrix}$$
, $Tr(A_0) = -1 < 0$.



Intuition: the noise mixes the two coordinates to yield stability on average.

Transport noise for parabolic equations

- M. Capiński's conjecture (around 1987):
 - In infinite dimensional case, stochastic transport noise could have a similar stabilizing effect for parabolic equations.
- Heat equation with transport noise:

$$df = \Delta f dt + \nu \sum_{k} \sigma_k \cdot \nabla f \circ dW_t^k,$$

where $\{\sigma_k\}_k$ are divergence free vector fields. Since

$$\langle \sigma_k \cdot \nabla f, g \rangle_{L^2} = -\langle f, \sigma_k \cdot \nabla g \rangle_{L^2},$$

transport noise behaves similarly as skew-symmetric matrices.

• By Stratonovich calculus,

$$\mathrm{d} \|f\|_{L^2}^2 = 2\langle f, \Delta f \rangle \, \mathrm{d} t + 2\nu \sum_k \langle f, \sigma_k \cdot \nabla f \rangle \circ \mathrm{d} W_t^k = -2 \|\nabla f\|_{L^2}^2 \, \mathrm{d} t.$$



Heuristic explanation



- Stirring the fluid produces small scale motions, which correspond to high frequencies in Fourier series.
- Higher frequencies correspond to smaller eigenvalues of Δ $(-\lambda_k \searrow -\infty)$, which yield stronger dissipation.
- Since transport noise mixes Fourier modes, we expect the top Lyapunov exponent

$$\lim_{\nu \to \infty} \lambda_{\nu} = -\infty.$$

Some results in deterministic setting

• Constantin-Kiselev-Ryzhik-Zlatoš (Ann. Math., 2008) considered heat equation on compact manifolds M:

$$\partial_t f^{\nu} = \Delta f^{\nu} + \nu \, b \cdot \nabla f^{\nu}, \quad f^{\nu}(0, x) = f_0(x),$$

where b is a Lipschitz divergence free vector field.

They studied the phenomenon of dissipation enhancement.

Theorem

If $b \cdot \nabla$ has no eigenfunctions in $H^1(M)$ (except trivial constant functions), then $\forall t > 0$ and $f_0 \in L^2(M)$, one has

$$\lim_{\nu \to \infty} \| f^{\nu}(t, \cdot) - \bar{f}_0 \|_{L^2} = 0,$$

where $\bar{f}_0 = \frac{1}{|M|} \int_M f_0 \, \mathrm{d}x$.



Some results in deterministic setting

- A. Zlatoš, Diffusion in fluid flow: dissipation enhancement by flows in 2D. Commun. Partial Differ. Equ. 35(3), 496–534 (2010)
- J. Bedrossian, M. Coti Zelati, Enhanced dissipation, hypoellipticity, and anomalous small noise inviscid limits in shear flows. *Arch. Rat. Mech. Anal.* 224(3), 1161–1204 (2017)
- Y. Feng, G. Iyer, Dissipation enhancement by mixing. *Nonlinearity* 32(5), 1810–1851 (2019)
- M. Coti Zelati, M.G. Delgadino, T.M. Elgindi, On the relation between enhanced dissipation timescales and mixing rates. Commun. Pure Appl. Math. 73(6), 1205–1244 (2020)
- J. Bedrossian, A. Blumenthal, S. Punshon-Smith, Almost-sure enhanced dissipation and uniform-in-diffusivity exponential mixing for advection -diffusion by stochastic Navier-Stokes. *Probab. Theory Related Fields* 179(3-4), 777-834 (2021)
- B. Gess, I. Yaroslavtsev, Stabilization by transport noise and enhanced dissipation in the Kraichnan model. arXiv:2104.03949 (2021)

Transport noise in inviscid models

• 2D Euler equation:

$$\begin{cases} \partial_t u + u \cdot \nabla u + \nabla p = 0 \\ \nabla \cdot u = 0 \end{cases} \longleftrightarrow \begin{cases} \partial_t \xi + u \cdot \nabla \xi = 0 \\ \xi = \nabla^{\perp} \cdot u = \partial_2 u_1 - \partial_1 u_2 \end{cases}$$

The enstrophy meas. (or white noise meas.) is formally invariant:

$$\mu(d\xi) = \frac{1}{Z} \exp\left(-\frac{1}{2} \|\xi\|_{L^2}^2\right) d\xi, \quad \text{supp}(\mu) = H^{-1-}(\mathbb{T}^2).$$

• Flandoli-Luo (AoP, 2020): white noise solutions of stoch. 2D Euler eqs

$$\begin{split} \mathrm{d}\xi^N + u^N \cdot \nabla \xi^N \, \mathrm{d}t &= \sqrt{\frac{2\nu}{\log N}} \sum_{k \in \mathbb{Z}_0^2, |k| \le N} \frac{1}{|k|} \sigma_k \cdot \nabla \xi^N \circ \mathrm{d}W_t^k \\ &= \nu \Delta \xi^N \, \mathrm{d}t + \sqrt{\frac{2\nu}{\log N}} \sum_{k \in \mathbb{Z}_0^2, |k| \le N} \frac{1}{|k|} \sigma_k \cdot \nabla \xi^N \, \mathrm{d}W_t^k \end{split}$$

converge to the unique stationary solu. of 2D NSEs with space-time white noise:

$$\partial_t \xi + u \cdot \nabla \xi = \nu \Delta \xi + \sqrt{2\nu} \nabla^{\perp} \cdot \eta \longleftrightarrow \partial_t u + u \cdot \nabla u + \nabla p = \nu \Delta u + \sqrt{2\nu} \eta,$$

where $\eta(t, x) = \sum_k \sigma_k(x) \dot{W}_t^k$, $\sigma_k(x) = \frac{k^{\perp}}{|k|} e_k(x)$.

Stoch. linear transport eqs

• Lucio Galeati (SPDE Anal Comp, 2020): stochastic linear transport eqs with L^2 -initial data f_0 :

$$df = b \cdot \nabla f \, dt + \sqrt{2\nu} \sum_{k \in \mathbb{Z}_0^2} \theta_k \sigma_k \cdot \nabla f \circ dW_t^k$$
$$= b \cdot \nabla f \, dt + \nu \Delta f \, dt + \sqrt{2\nu} \sum_k \theta_k \sigma_k \cdot \nabla f \, dW_t^k,$$

where $\nu > 0$ is noise intensity, $\{\theta_k\}_k \in \ell^2(\mathbb{Z}_0^2)$, $\theta_k = \theta_l$ for all |k| = |l|. $\forall t \leq T$, it holds $||f_t||_{L^2} \leq C_T ||f_0||_{L^2}$ (independent of ν and $\theta \in \ell^2$).

• Taking $\{\theta^N\}_{N\geq 1}\subset \ell^2(\mathbb{Z}_0^2)$ satisfying $\|\theta^N\|_{\ell^2}=1$ and $\lim_{N\to\infty}\|\theta^N\|_{\ell^\infty}=0$, Lucio Galeati proved

$$df^{N} = b \cdot \nabla f^{N} dt + \nu \Delta f^{N} dt + \sqrt{2\nu} \sum_{k} \theta_{k}^{N} \sigma_{k} \cdot \nabla f^{N} dW_{t}^{k}$$

converge to the deterministic parabolic eq:

$$\partial_t f = b \cdot \nabla f + \nu \Delta f, \quad f|_{t=0} = f_0.$$

• If $\theta_k^N = \frac{1}{Z_N} \frac{\mathbf{1}_{\{N \le |k| \le 2N\}}}{|k|^a} (k \in \mathbb{Z}_0^2)$, then $Z_N \sim \frac{1}{N^{a-1}}$ and $\|\theta^N\|_{\ell^\infty} \sim \frac{1}{N}$.

Why the martingale part vanishes?

Recall the martingale part

$$\mathrm{d} M^N_t = \sum_{k \in \mathbb{Z}^2_0} \theta^N_k \sigma_k \cdot \nabla f^N_t \, \mathrm{d} W^k_t.$$

For any $\phi \in C^1(\mathbb{T}^2)$, we have

$$\langle M_t^N, \phi \rangle = -\sum_k \theta_k^N \int_0^t \langle f_s^N, \sigma_k \cdot \nabla \phi \rangle dW_s^k.$$

$$\implies \mathbb{E} \langle M_t^N, \phi \rangle^2 = \sum_k (\theta_k^N)^2 \mathbb{E} \int_0^t \langle f_s^N, \sigma_k \cdot \nabla \phi \rangle^2 ds$$

$$\leq \|\theta^N\|_{\ell^\infty}^2 \mathbb{E} \int_0^t \sum_k \langle f_s^N \nabla \phi, \sigma_k \rangle^2 ds$$

$$\leq \|\theta^N\|_{\ell^\infty}^2 \mathbb{E} \int_0^t \|f_s^N \nabla \phi\|_{L^2}^2 ds$$

$$\leq \|\theta^N\|_{\ell^\infty}^2 \|\nabla \phi\|_{L^\infty}^2 C_T \|f_0\|_{L^2}^2 \to 0.$$

Scaling limit of stoch 2D Euler eq

• Flandoli-Galeati-Luo (*J Evol Equ*, 2021): stochastic 2D Euler equations in vorticity form, with L²-initial data:

$$d\xi + u \cdot \nabla \xi dt = \sqrt{2\nu} \sum_{k} \theta_{k} \sigma_{k} \cdot \nabla \xi \circ dW_{t}^{k}$$
$$= \nu \Delta \xi dt + \sqrt{2\nu} \sum_{k} \theta_{k} \sigma_{k} \cdot \nabla \xi dW_{t}^{k}.$$

 \exists weak solution $\{\xi_t\}_{t\in[0,T]}$ such that \mathbb{P} -a.s. $\|\xi_t\|_{L^2} \leq \|\xi_0\|_{L^2}$.

• Taking $\{\theta^N\}_N \subset \ell^2(\mathbb{Z}_0^2)$ s.t. $\|\theta^N\|_{\ell^2} = 1$ and $\lim_{N \to \infty} \|\theta^N\|_{\ell^\infty} = 0$, we proved $\xi^N \stackrel{w}{\rightharpoonup} \xi$, the unique solution of the deterministic 2D NSE in vorticity form:

$$\begin{cases} \partial_t \xi + u \cdot \nabla \xi = \nu \Delta \xi, \\ u = K * \xi \end{cases} \longleftrightarrow \begin{cases} \partial_t u + u \cdot \nabla u + \nabla p = \nu \Delta u, \\ \nabla \cdot u = 0. \end{cases}$$

• The proof is based on compactness arguments, hence there is no convergence rate.

Related works: suppression of blow-up

• 3D NSEs on \mathbb{T}^3 .

$$\begin{cases} \partial_t u + u \cdot \nabla u + \nabla p = \Delta u \\ \nabla \cdot u = 0 \end{cases} \longleftrightarrow \begin{cases} \partial_t \xi + u \cdot \nabla \xi - \xi \cdot \nabla u = \Delta \xi \\ \xi = \nabla \times u \end{cases}$$

 $u, \, \xi : \mathbb{R}_+ \times \mathbb{T}^3 \to \mathbb{R}^3.$

• Flandoli-Luo (PTRF, 2021): Stoch 3D NSEs.

$$d\xi + (u \cdot \nabla \xi - \xi \cdot \nabla u) dt = \Delta \xi dt + \sqrt{2\nu} \sum_{k \in \mathbb{Z}_0^3} \sum_{i=1}^2 \theta_k \Pi(\sigma_{k,i} \cdot \nabla \xi) \circ dW_t^{k,i}.$$

 $\forall R > 0, \exists \text{ big } \nu \text{ and } \theta \in \ell^2(\mathbb{Z}_0^3), \text{ such that } \forall \|\xi_0\|_{L^2} \leq R, \text{ global solution exists with large probability.}$

 \bullet Flandoli-Galeati-Luo (Comm. PDE, 2021): general nonlinear eqs on $\mathbb{T}^d.$

$$df = -(-\Delta)^{\alpha} f dt + \Phi(f) dt + \sqrt{2\nu} \sum_{k \in \mathbb{Z}_0^d} \sum_{i=1}^{d-1} \theta_k \sigma_{k,i} \cdot \nabla f \circ dW_t^{k,i},$$

where $\alpha \geq 1$ and $\Phi: H^{\alpha} \to H^{-\alpha}$ is some nonlinear mapping.



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Stochastic heat equation

Heat equation with transport noise on \mathbb{T}^2 :

$$\partial_t f = \kappa \Delta f + \dot{W}_t \cdot \nabla f, \quad W_t(x) = \sqrt{2\nu} \sum_{k \in \mathbb{Z}_0^2} \theta_k \sigma_k(x) W_t^k;$$

- $\kappa > 0$: molecular diffusivity; $\nu > 0$: noise intensity.
- $\theta = (\theta_k)_{k \in \mathbb{Z}_0^2} \in \ell^2(\mathbb{Z}_0^2), \|\theta\|_{\ell^2} = 1, \theta$ is radially symmetric.
- $\bullet \ \sigma_k(x) = \frac{k^{\perp}}{|k|} e_k(x), \ k \in \mathbb{Z}_0^2.$

Stratonovich to Itô:

$$df = \kappa \Delta f dt + \sqrt{2\nu} \sum_{k} \theta_{k} \sigma_{k} \cdot \nabla f \circ dW_{t}^{k}$$

$$= \kappa \Delta f dt + \nu \sum_{k} \theta_{k}^{2} \sigma_{k} \cdot \nabla (\sigma_{k} \cdot \nabla f) dt + \sqrt{2\nu} \sum_{k} \theta_{k} \sigma_{k} \cdot \nabla f dW_{t}^{k}$$

$$= (\kappa + \nu) \Delta f dt + \sqrt{2\nu} \sum_{k} \theta_{k} \sigma_{k} \cdot \nabla f dW_{t}^{k}.$$

Dissipation enhancement

- Though there is $\nu\Delta$, it does not mean dissipation has already been enhanced.
- Indeed, by Itô's formula,

$$\begin{split} \mathrm{d}\|f\|_{L^{2}}^{2} &= 2\kappa \langle f, \Delta f \rangle \, \mathrm{d}t + 2\sqrt{2\nu} \sum_{k} \theta_{k} \langle f, \sigma_{k} \cdot \nabla f \rangle \circ \mathrm{d}W_{t}^{k} \\ &= -2\kappa \|\nabla f\|_{L^{2}}^{2} \, \mathrm{d}t \leq -8\kappa \pi^{2} \|f\|_{L^{2}}^{2} \, \mathrm{d}t \\ \Longrightarrow \|f_{t}\|_{L^{2}} \leq \|f_{0}\|_{L^{2}} \, e^{-4\kappa \pi^{2}t} \quad \text{(independent of } \nu \text{ and } \theta) \end{split}$$

Theorem (Flandoli-Galeati-Luo, arXiv:2104.01740)

 $\forall \lambda > 0, \exists \nu > 0 \text{ and } \theta \in \ell^2(\mathbb{Z}_0^2) \text{ with the following property:}$

 $\forall f_0 \in L^2(\mathbb{T}^d)$ with zero mean, \exists a random constant C > 0 such that, for the solution f_t with initial condition f_0 , it holds

$$\mathbb{P}$$
-a.s. $||f_t||_{L^2} \le Ce^{-\lambda t}$ for all $t \ge 0$.

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Stoch heat eqn and mild formulation

$$df = \kappa \Delta f dt + \sqrt{2\nu} \sum_{k} \theta_{k} \sigma_{k} \cdot \nabla f \circ dW_{t}^{k}$$
$$= (\kappa + \nu) \Delta f dt + \sqrt{2\nu} \sum_{k} \theta_{k} \sigma_{k} \cdot \nabla f dW_{t}^{k}.$$

• Energy balance: $\forall 0 \le s < t$,

$$||f_t||_{L^2}^2 + 2\kappa \int_s^t ||\nabla f_r||_{L^2}^2 dr = ||f_s||_{L^2}^2,$$

in particular, $t \to ||f_t||_{L^2}^2$ is decreasing.

• Mild formulation: $\forall 0 \le s < t$,

$$f_t = e^{(\kappa+\nu)(t-s)\Delta} f_s + \underbrace{\sqrt{2\nu} \sum_k \theta_k \int_s^t e^{(\kappa+\nu)(t-r)\Delta} (\sigma_k \cdot \nabla f_r) \, dW_r^k}_{Z_{s,t}}.$$



Key lemma

Lemma

 $\exists \delta \in (0,1) \text{ such that for any } n \geq 0,$

$$\mathbb{E} \|f_{n+1}\|_{L^2}^2 \le \delta \, \mathbb{E} \|f_n\|_{L^2}^2,$$

where, for some $0 < \alpha < 1 < \beta$,

$$\delta \le C_{\alpha,\beta} \left[(\kappa + \nu)^{-1} + \nu^{\frac{4\beta - \alpha(2\beta + d)}{4(\alpha + \beta)}} \kappa^{-\frac{\beta}{\alpha + \beta}} \|\theta\|_{\ell^{\infty}}^{\frac{2\alpha}{\alpha + \beta}} \right].$$

In particular, δ can be as small as possible by first taking ν big, and then letting $\|\theta\|_{\ell^{\infty}}$ small enough.

Idea of proof: using monotonicity $t \to ||f_t||_{L^2}^2$ and mild formulation $f_t = e^{(\kappa + \nu)(t-s)\Delta} f_s + Z_{s,t}$,

$$||f_{n+1}||_{L^{2}}^{2} \leq \int_{n}^{n+1} ||f_{t}||_{L^{2}}^{2} dt \leq 2 \int_{n}^{n+1} ||e^{(\kappa+\nu)(t-n)\Delta} f_{n}||_{L^{2}}^{2} dt + \int_{n}^{n+1} ||Z_{n,t}||_{L^{2}}^{2} dt$$
$$\leq 2 \int_{n}^{n+1} e^{8\pi^{2}(\kappa+\nu)(t-n)} ||f_{n}||_{L^{2}}^{2} dt + \int_{n}^{n+1} ||Z_{n,t}||_{L^{2}}^{2} dt.$$

Proof of the main theorem

Theorem (Flandoli-Galeati-Luo, arXiv:2104.01740)

 $\forall \lambda > 0, \exists \nu > 0 \text{ and } \theta \in \ell^2(\mathbb{Z}_0^2) \text{ with the following property:}$

 $\forall f_0 \in L^2(\mathbb{T}^d)$ with zero mean, \exists a random constant C > 0 such that, for the solution f_t with initial condition f_0 , it holds

$$\mathbb{P}$$
-a.s. $||f_t||_{L^2} \le Ce^{-\lambda t}$ for all $t \ge 0$.

Proof. The above lemma implies $\exists \delta \in (0,1)$ such that for any $n \geq 1$,

$$\mathbb{E}\left[\sup_{t\in[n,n+1]} \|f_t\|_{L^2}^2\right] = \mathbb{E}\|f_n\|_{L^2}^2 \le \delta \,\mathbb{E}\|f_{n-1}\|_{L^2}^2 \le \cdots$$
$$\le \delta^n \|f_0\|_{L^2}^2 = e^{-\lambda_0 n} \|f_0\|_{L^2}^2.$$

Choosing ν big and $\|\theta\|_{\ell^{\infty}}$ small, we can make λ_0 as big as we want.

Proof of the main theorem

Fix any $\lambda \in (0, \lambda_0)$, define the event

$$A_n = \left\{ \omega \in \Omega : \sup_{t \in [n, n+1]} \|f_t(\omega)\|_{L^2}^2 > e^{-\lambda n} \right\}, \quad n \ge 1.$$

Then by Chebyshev's inequality,

$$\sum_{n} \mathbb{P}(A_n) \le \sum_{n} e^{\lambda n} \mathbb{E}\left[\sup_{t \in [n, n+1]} \|f_t\|_{L^2}^2\right] \le \sum_{n} e^{(\lambda - \lambda_0)n} \|f_0\|_{L^2}^2 < +\infty.$$

Borel-Cantelli lemma implies, for \mathbb{P} -a.s. ω , $\exists n(\omega) \in \mathbb{N}$ such that

$$\sup_{t \in [n, n+1]} \|f_t(\omega)\|_{L^2}^2 \le e^{-\lambda n}, \quad \forall n \ge n(\omega).$$

Thus \exists a random constant C > 0 such that

$$||f_t(\omega)||_{L^2}^2 \le Ce^{-\lambda t}, \quad t \ge 0.$$



谢谢大家!